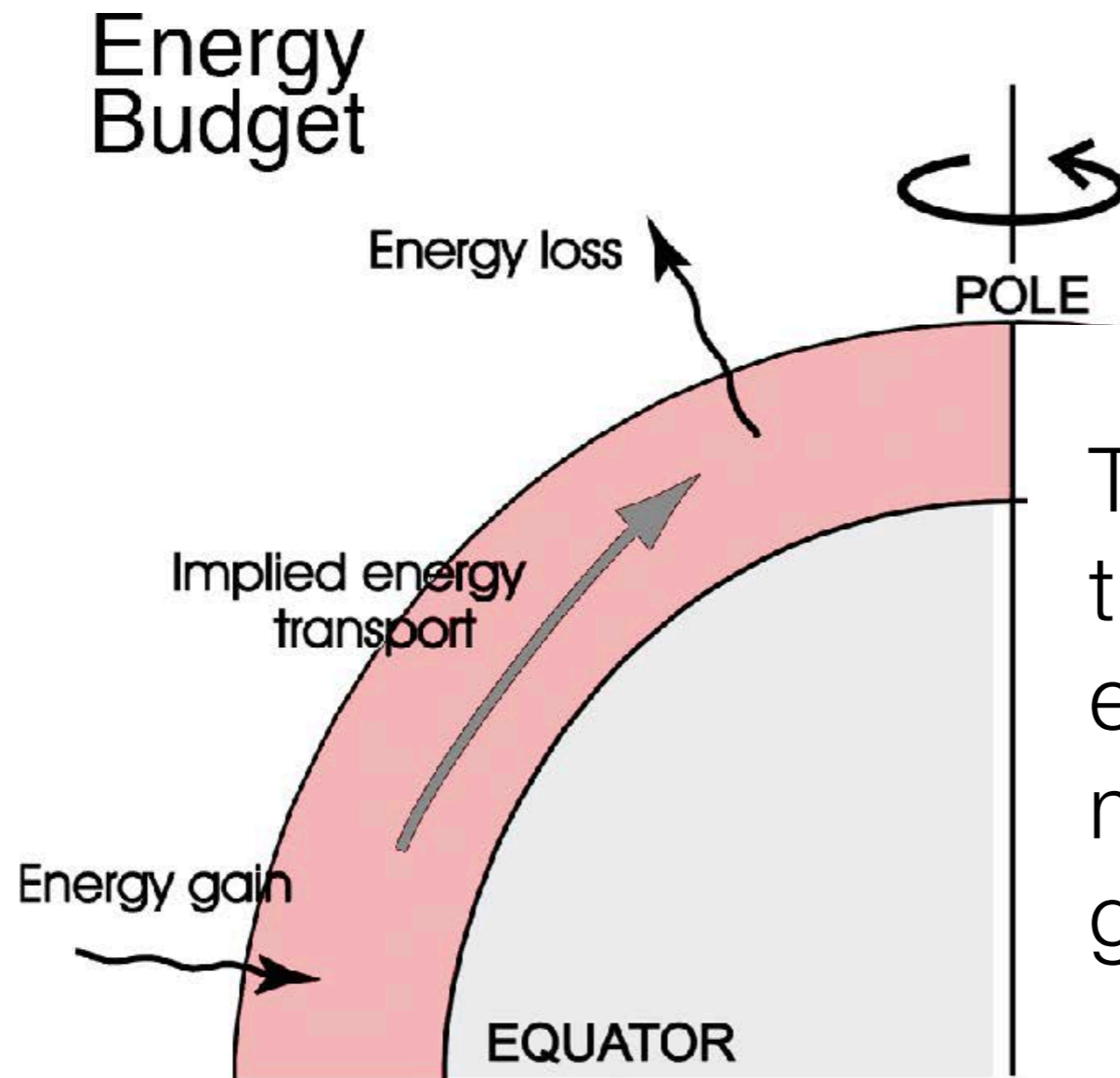


# Atmosphere: #5

## General circulation

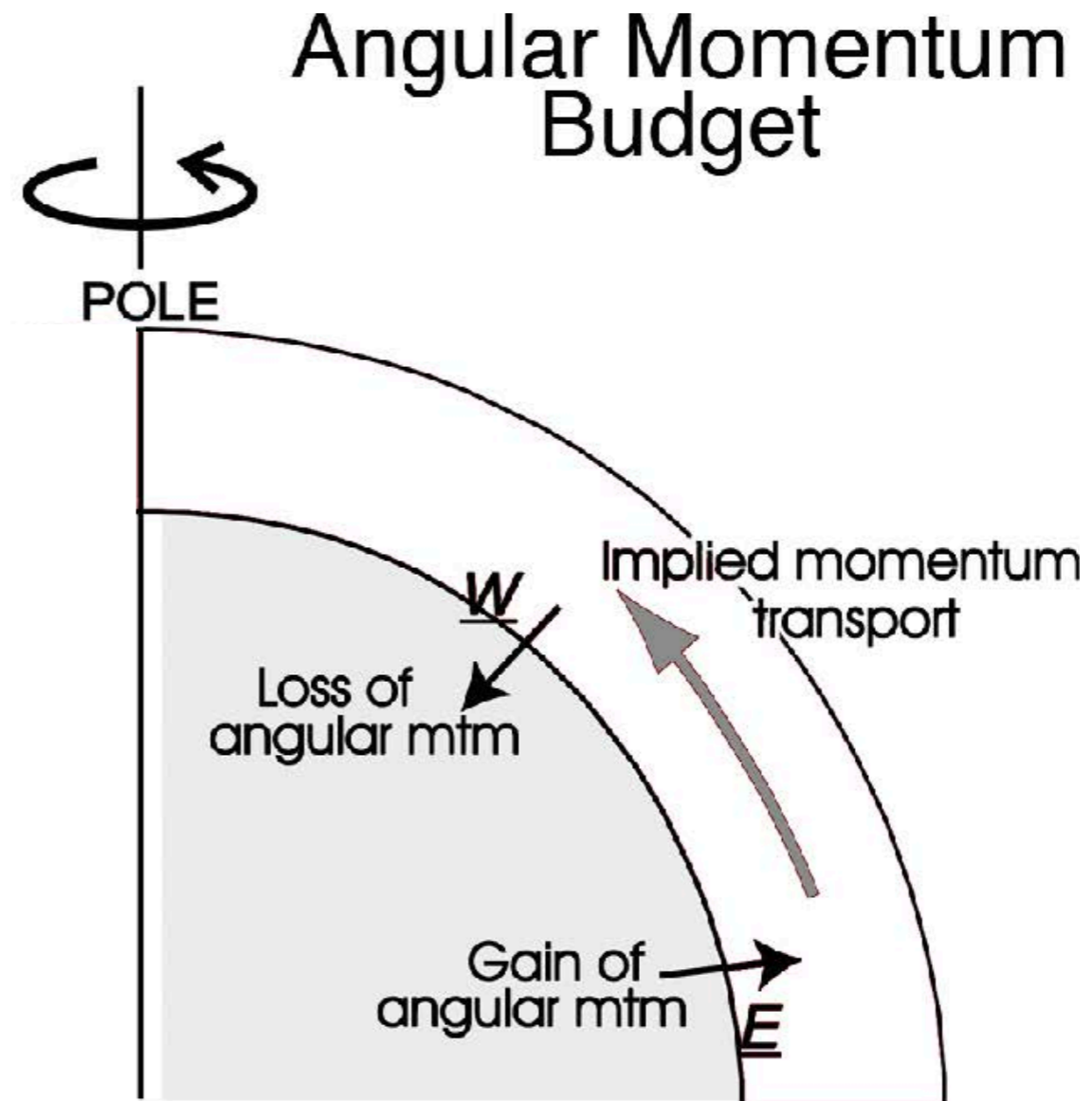
# General circulation of the atmosphere



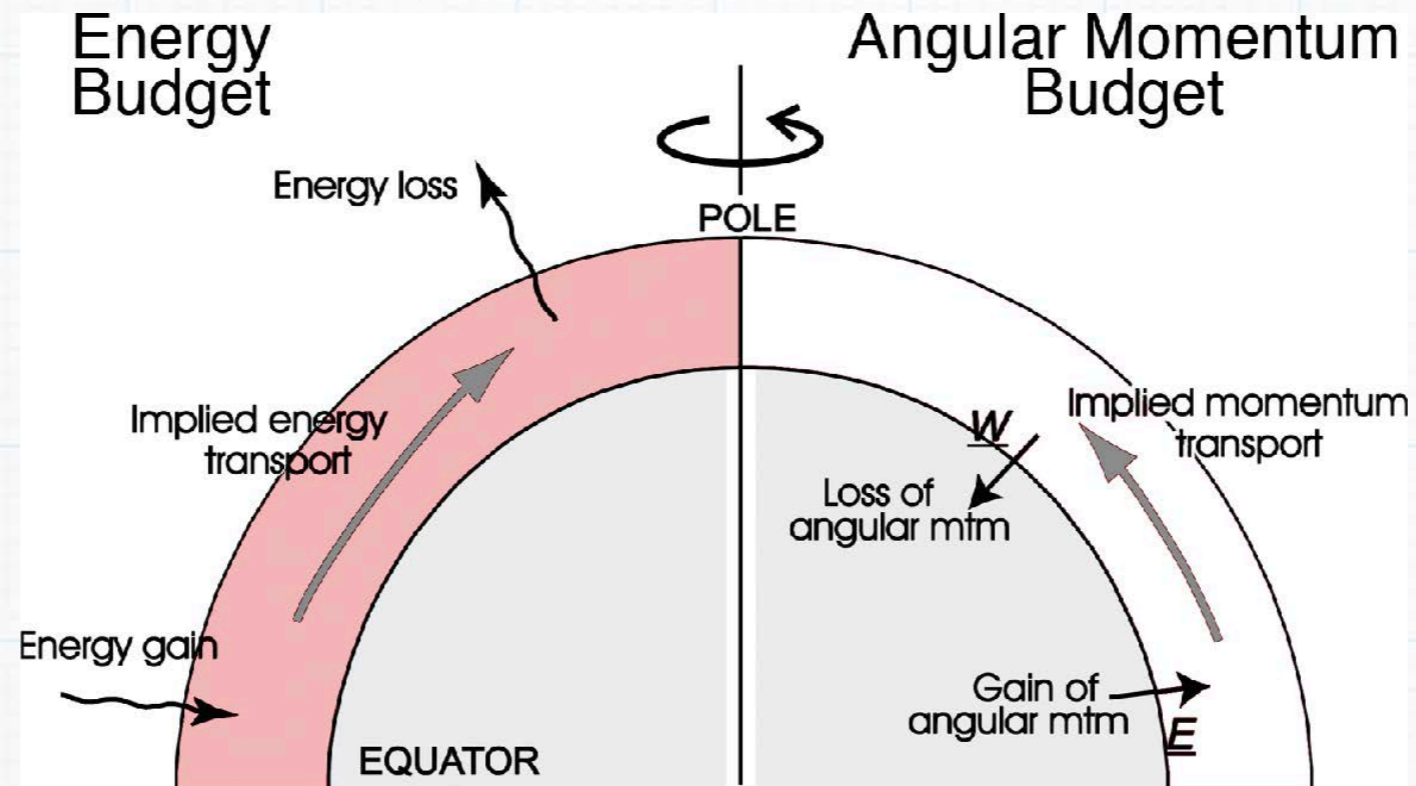
The atmosphere has to transport energy from equator to pole to maintain the temperature gradient

# General circulation of the atmosphere

The atmosphere has to transport westerly angular momentum from low to middle latitude.



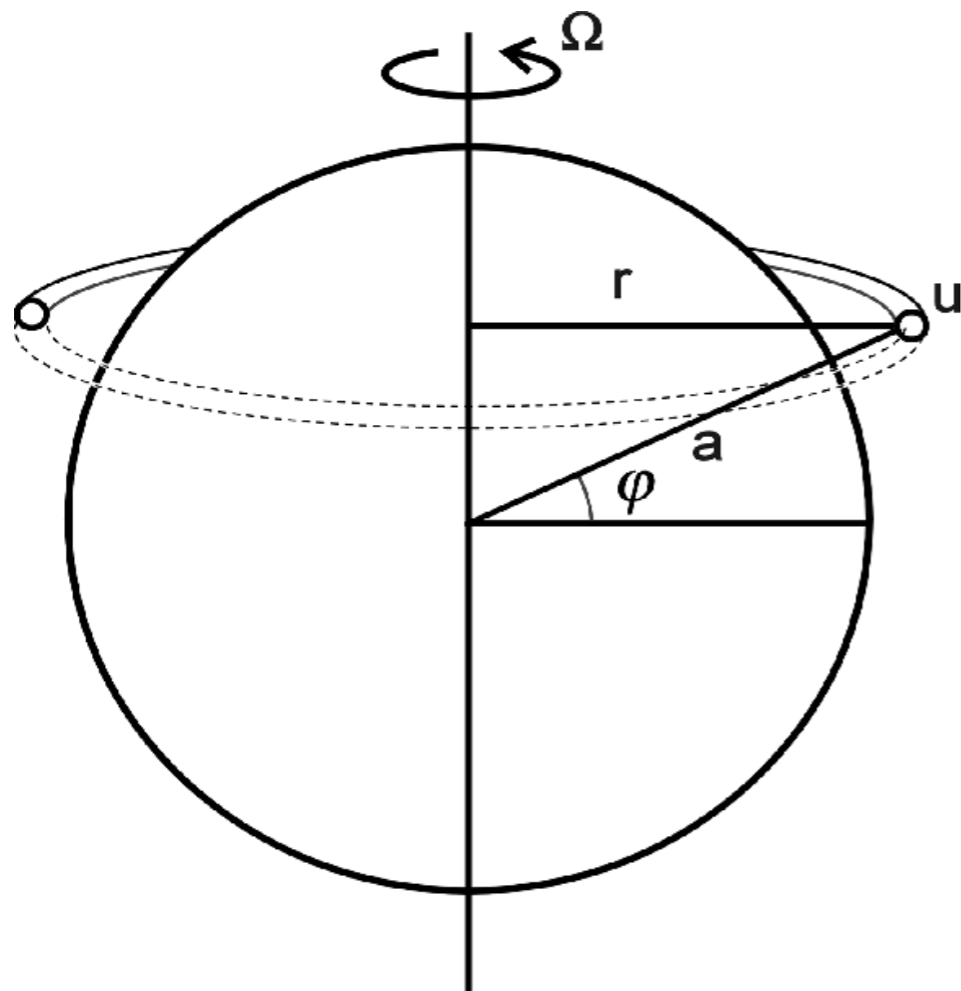
# General circulation of the atmosphere



- In the upper troposphere, we know that the dominant flow is the west-to-east, which cannot explain the equator-to-pole transport of heat and angular momentum.
- This is why the overturning circulation becomes important.
- But the meridional overturning circulation does not extend all the way to the pole as in the figure. Why?

# Mechanistic view of the circulation: tropics

- For simplicity, let's assume homogeneous surface with no seasonality.
- So we focus on the temperature gradient in latitudinal direction.



Absolute angular momentum

$$\begin{aligned} A &= \Omega r^2 + ur \\ &= \Omega a^2 \cos^2 \phi + ua \cos \phi \end{aligned}$$

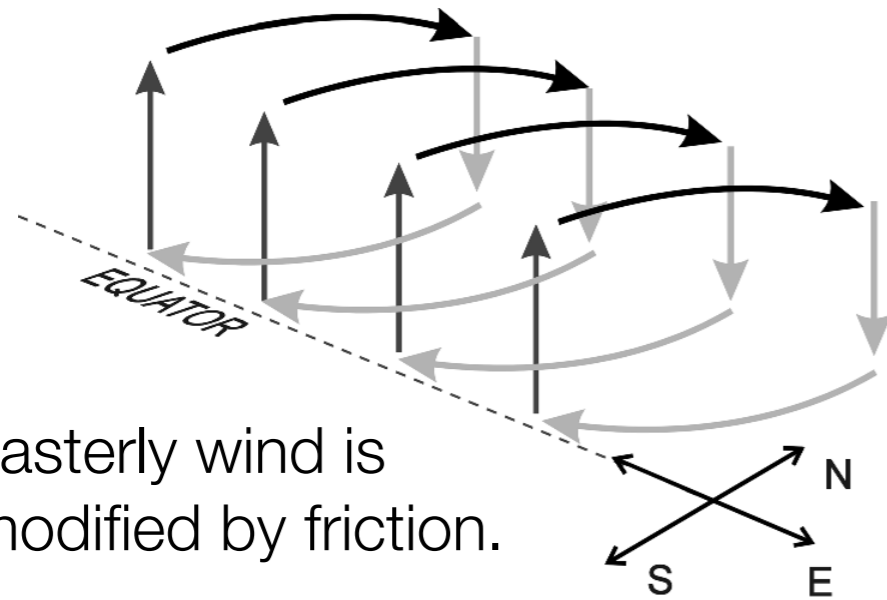
If we start  $u=0$  at the equator,  $u$  grows as we go to the north.

# Mechanistic view of the circulation: tropics

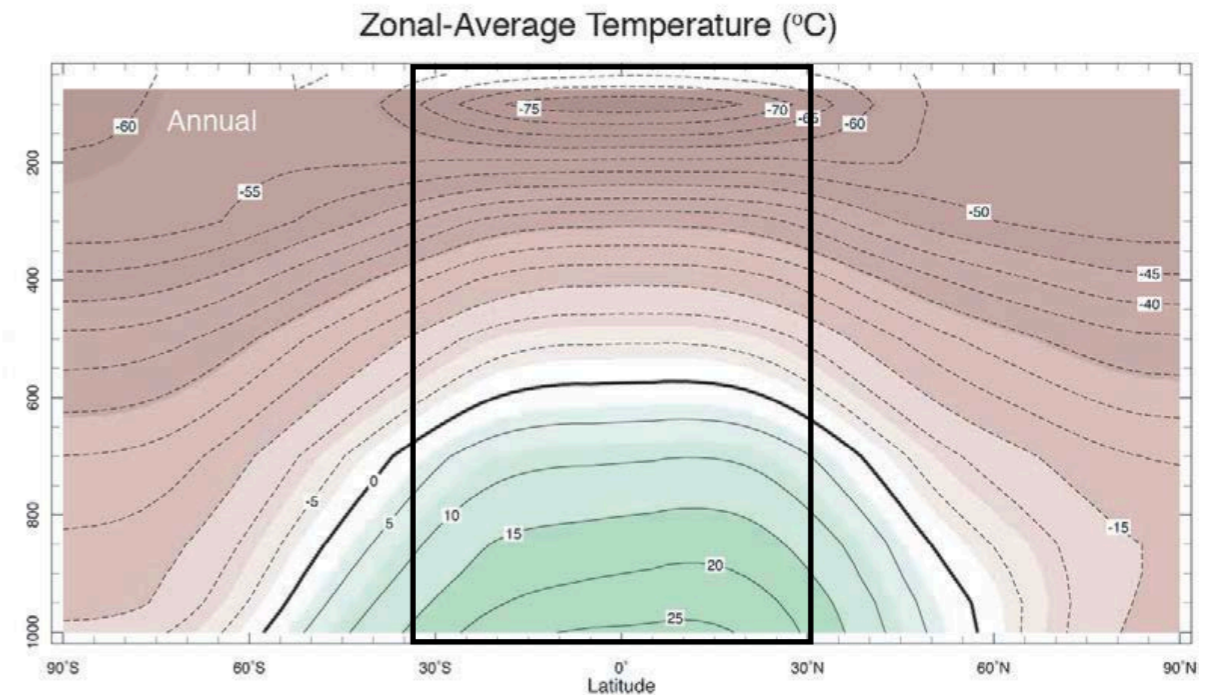
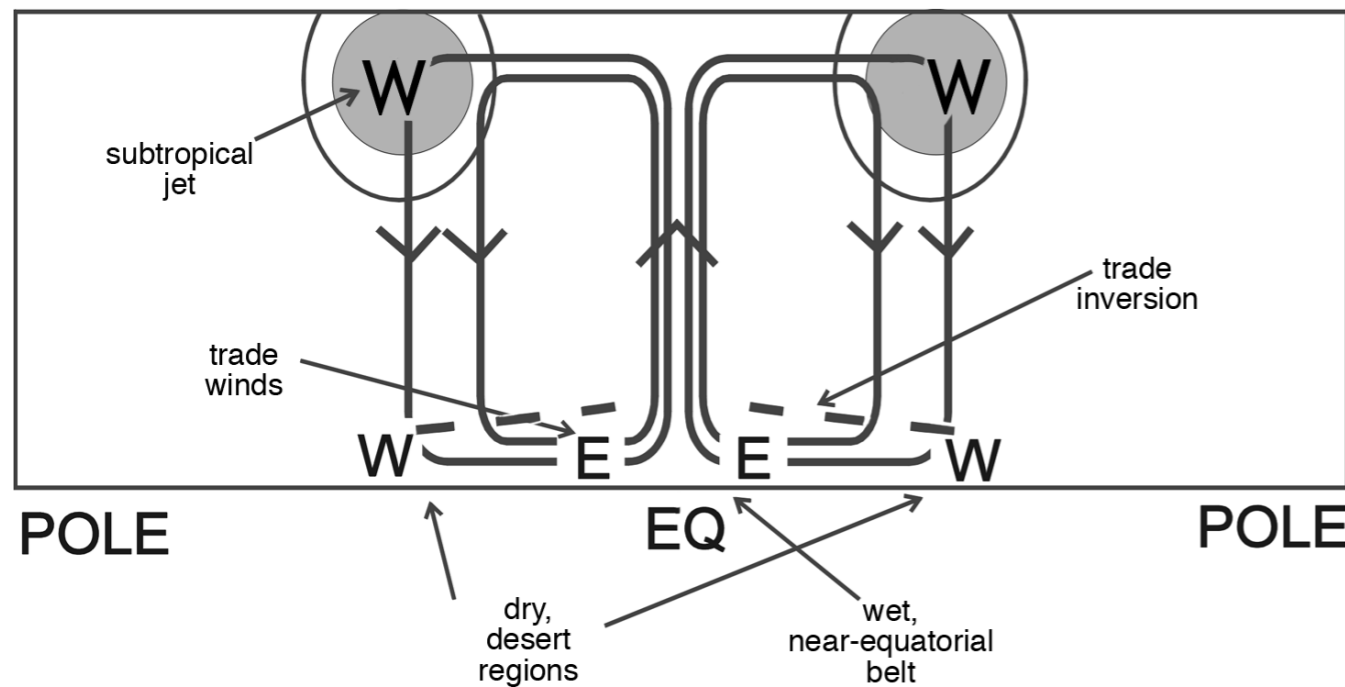
westerly wind

Subtropical jet is driven in large part by the advection of angular momentum by the Hadley cell.

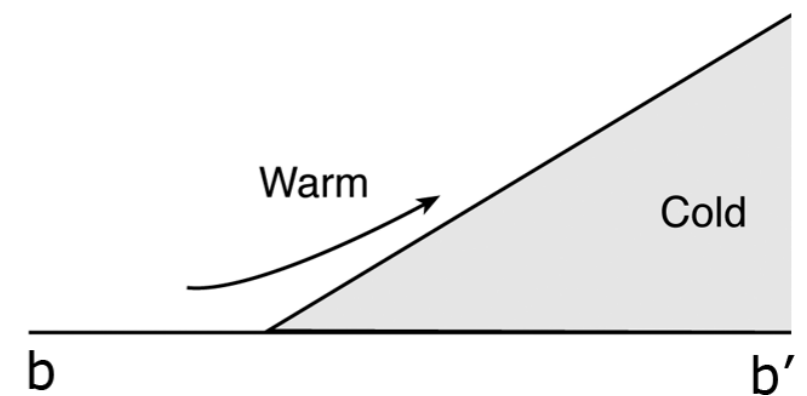
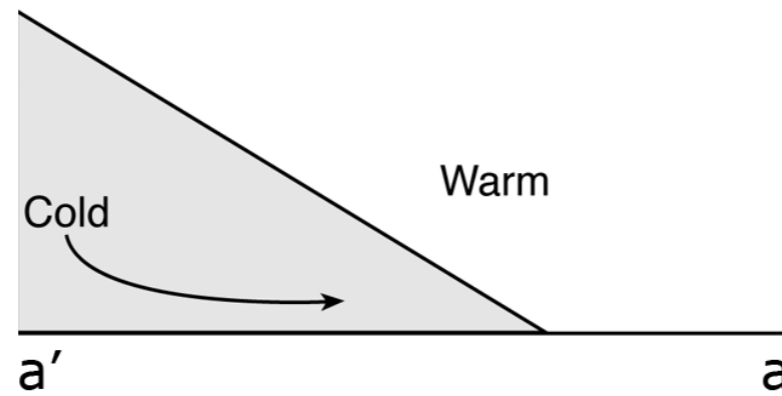
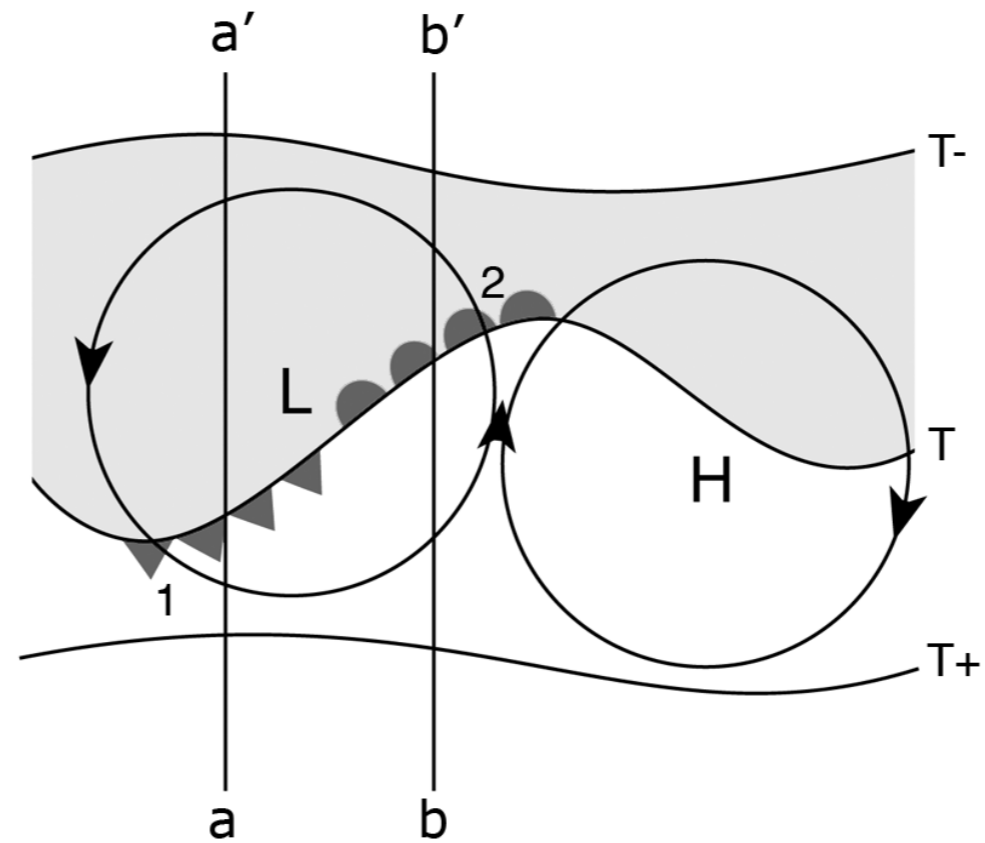
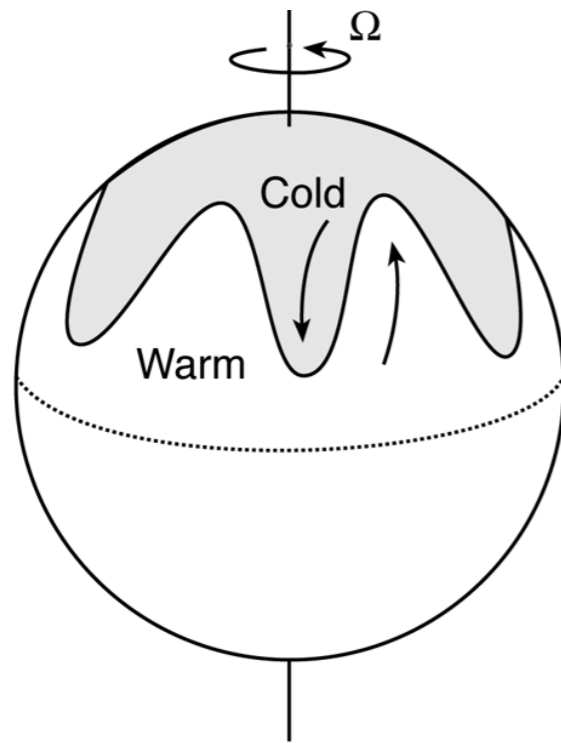
easterly wind is modified by friction.



Small T gradient reflects how effective Hadley cell is in transporting the heat.



# Mechanistic view of the circulation: extratropics



# Energetics of the thermal wind equation

Q: Why do we have this energetic flows?

A: Available potential energy that can be released by a redistribution of mass of the system



# Energetics of the thermal wind equation

- Let's consider an incompressible fluid, like water, for simplicity.
- A potential energy (PE) of a fluid parcel of volume  $dV = dx dy dz$  and density  $\rho$  would be  $gz\rho dV$ :
- The total potential energy is then

$$PE = g \int z\rho dV = g \int \rho dV \frac{\int z\rho dV}{\int \rho dV} = gM \langle z \rangle$$

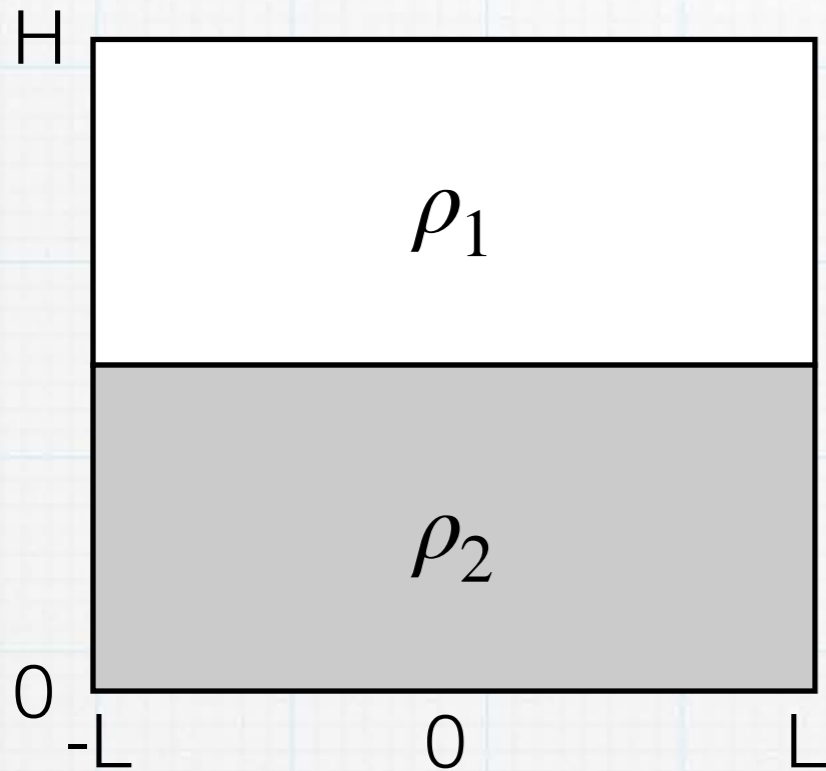
The height of the center of mass



# Energetics of the thermal wind equation

- Energy can be released and converted to kinetic energy only if some rearrangement of the fluid results in a lower total potential energy

Example:



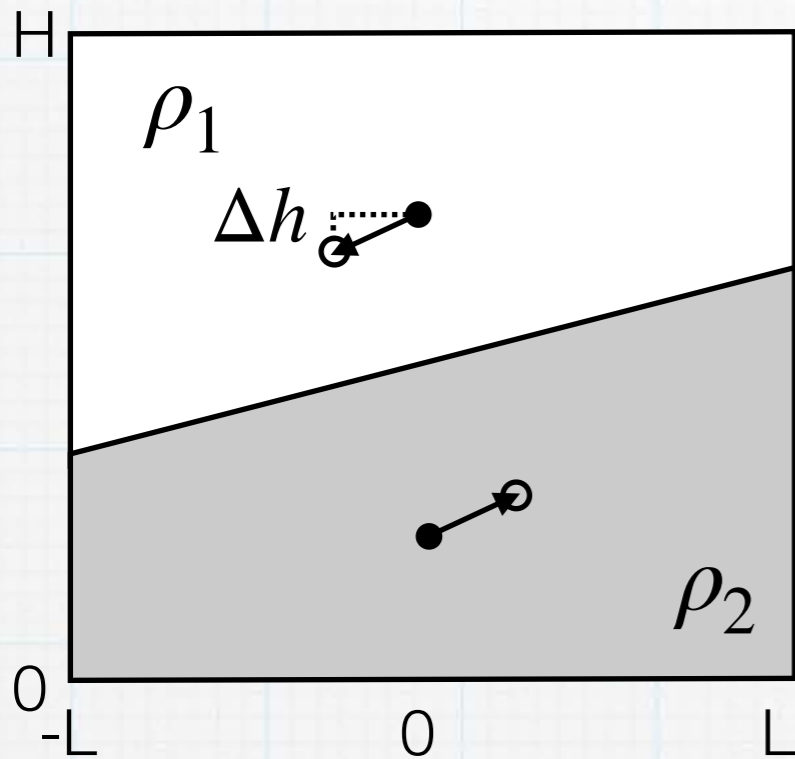
$$PE_1 = \rho_1 g \frac{3}{4} H$$

$$PE_2 = \rho_2 g \frac{1}{4} H$$

$$PE_A = gH \left( \frac{3\rho_1 + \rho_2}{4} \right)$$

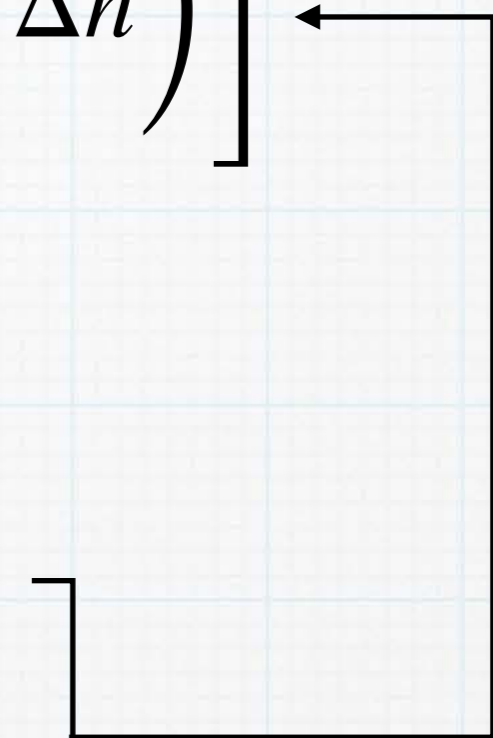
# Energetics of the thermal wind equation

$$PE_B = g \left[ \rho_1 \left( \frac{3}{4}H - \Delta h \right) + \rho_2 \left( \frac{1}{4}H + \Delta h \right) \right]$$



$$PE_1 = \rho_1 g \left( \frac{3}{4}H - \Delta h \right)$$

$$PE_2 = \rho_2 g \left( \frac{1}{4}H + \Delta h \right)$$



# Energetics of the thermal wind equation

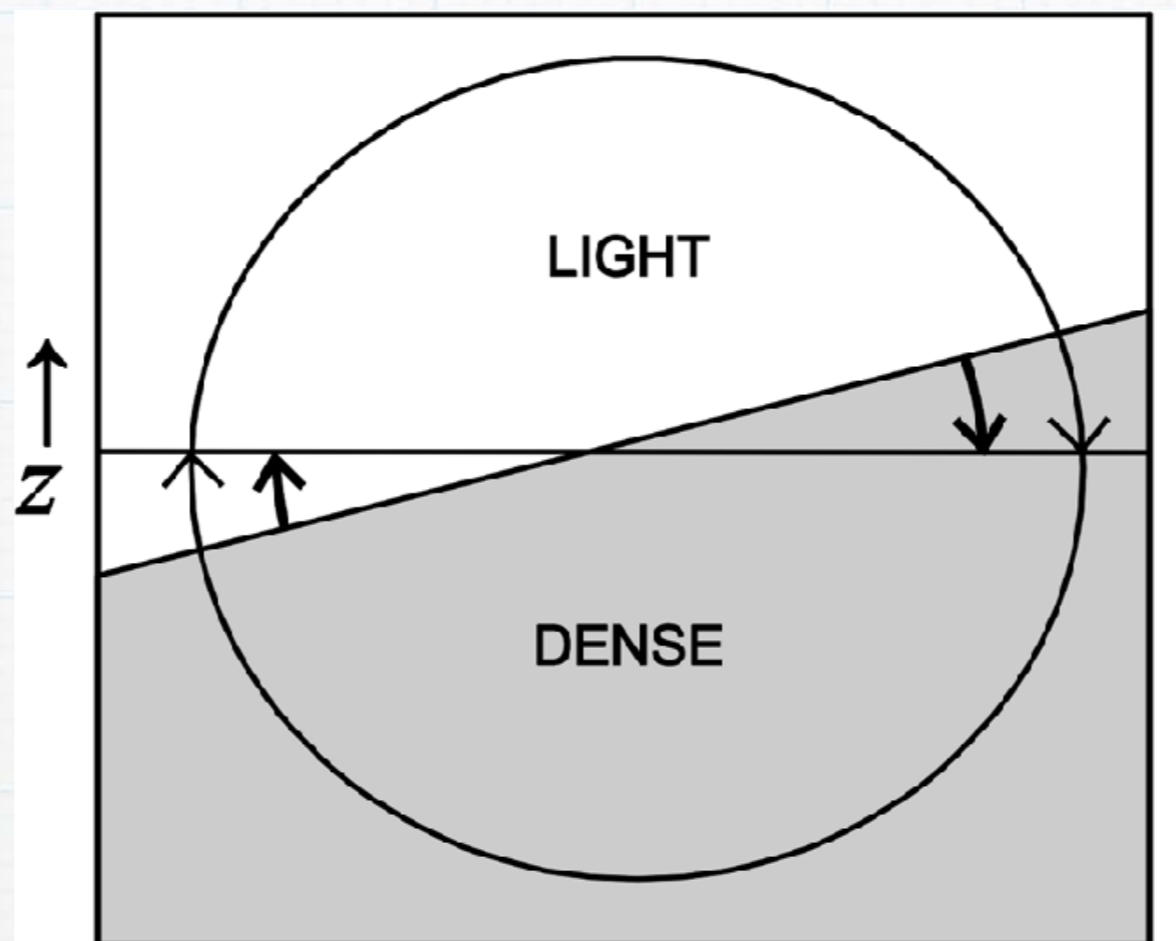
- Which fluid has higher potential energy?

$$PE_A - PE_B = g \left( \rho_1 \frac{3}{4} H + \rho_2 \frac{1}{4} H - \rho_1 \frac{3}{4} H + \Delta h \rho_1 - \rho_2 \frac{1}{4} H - \Delta h \rho_2 \right)$$
$$= g \Delta h (\rho_1 - \rho_2) < 0$$

- The case B has higher potential energy by  $g \Delta h (\rho_1 - \rho_2)$
- This is available potential energy (APE).
- We can expect higher available potential energy when the interface has greater tilt.

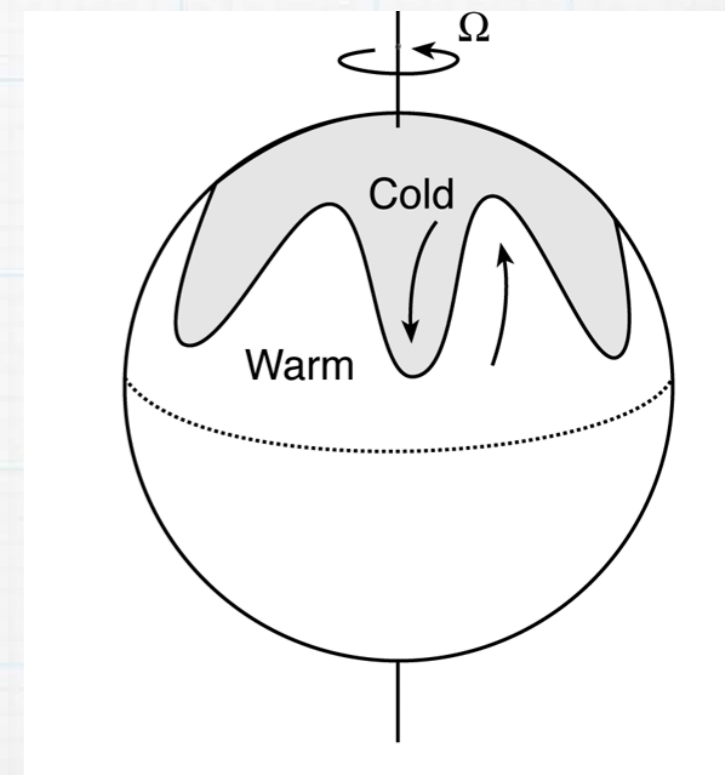
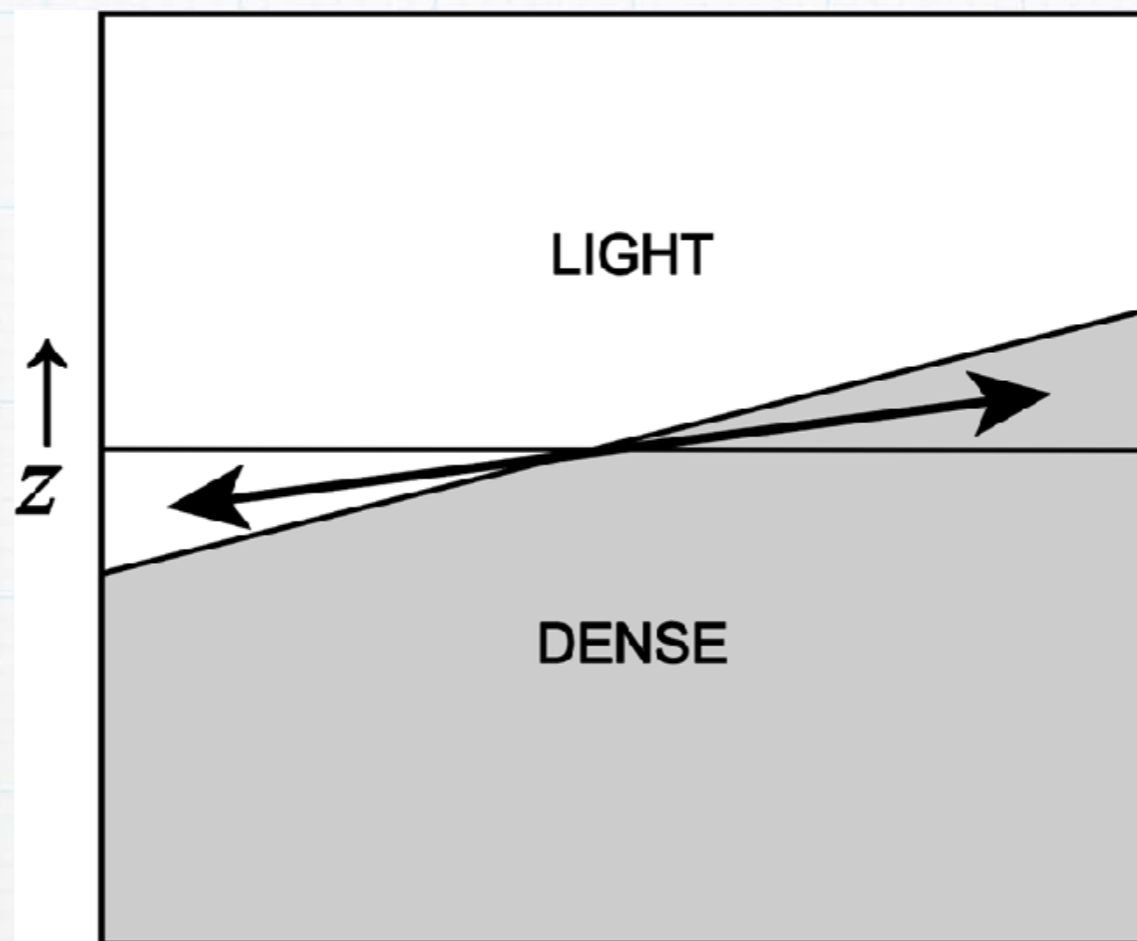
# Energetics of the thermal wind equation

- Release of available potential energy
  - In a non-rotating fluid



# Energetics of the thermal wind equation

- Release of available potential energy
  - In a rotating fluid, the tilted slope can be maintained in thermal wind balance

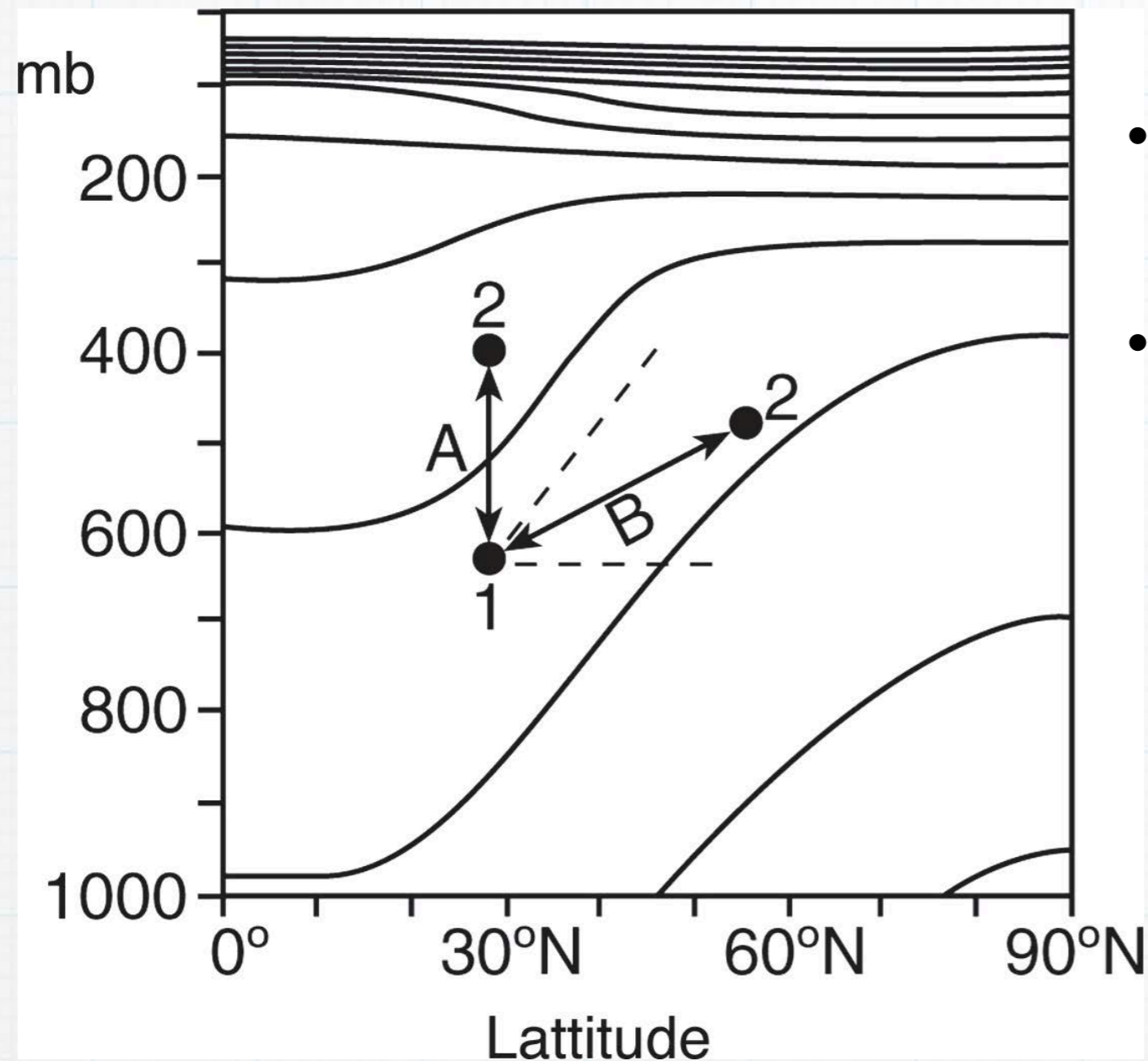


# Energetics in a compressible atmosphere

- The air is compressible → we need to consider internal energy.
  - Internal energy goes up when compressed
  - Internal energy goes down when expanded
- The air can contain moisture → we need to consider latent heat when condensation occurs.
- Total energy of the atmosphere = potential energy + kinetic energy + internal energy + latent heat content

# Energetics in a compressible atmosphere

Potential temperature (increasing with height)



- Moving from 1 to 2 along A : needs energy
- Moving from 1 to 2 along B : release energy → can excite eddies



# 1. Energy transport

- Total energy of the atmosphere = internal energy + potential energy + latent heat content + kinetic energy

$$E = c_p T + gz + Lq + \frac{1}{2} \mathbf{u} \cdot \mathbf{u}$$

- Energy transport by the atmosphere across the unit area =  $\rho v E dA$

- Total meridional energy transport =  $\iint \rho v E dx dz$

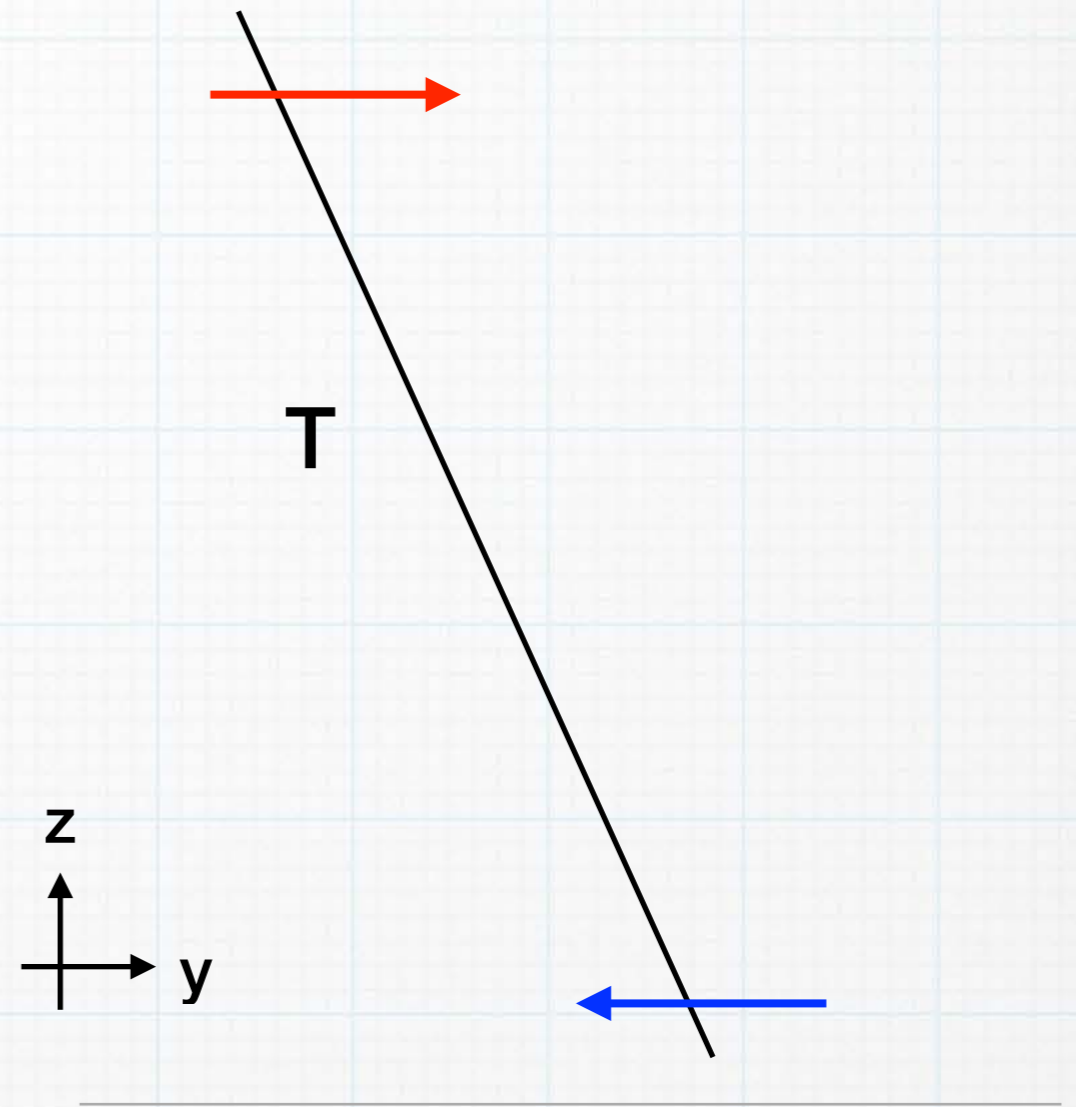
# 1. Energy transport, tropics

- The internal energy

$$\int_0^{\infty} \rho v c_p T dz < 0$$

Equatorward heat transport

The Hadley circulation carries heat toward the hot equator from the cooler subtropics!



# 1. Energy transport, tropics

- The internal energy + potential energy

$$\int_0^{\infty} \rho v (c_p T + gz) dz = c_p \int_0^{\infty} \rho v \left( T + \frac{g}{c_p} z \right) dz$$
$$= c_p \int_0^{\infty} \rho v \left( T - \frac{dT}{dz} \Big|_{dry} z \right) dz > 0$$

The atmosphere is stable in dry adiabatic process,  
which makes this term positive.

The Hadley circulation carries (heat+potential) energy poleward.

# 1. Energy transport, tropics

- Upper branch has far less moisture than lower branch of the Hadley cell.
- The net latent heat transport by the Hadley cell is equatorward.
- It turned out that poleward (heat+potential) energy transport and equatorward latent heat energy transport are in opposite sign with similar magnitude.
- The kinetic energy has negligible contribution to the total energy.
- In the net, then, the annually averaged energy flux by the Hadley cell is (weakly) poleward.

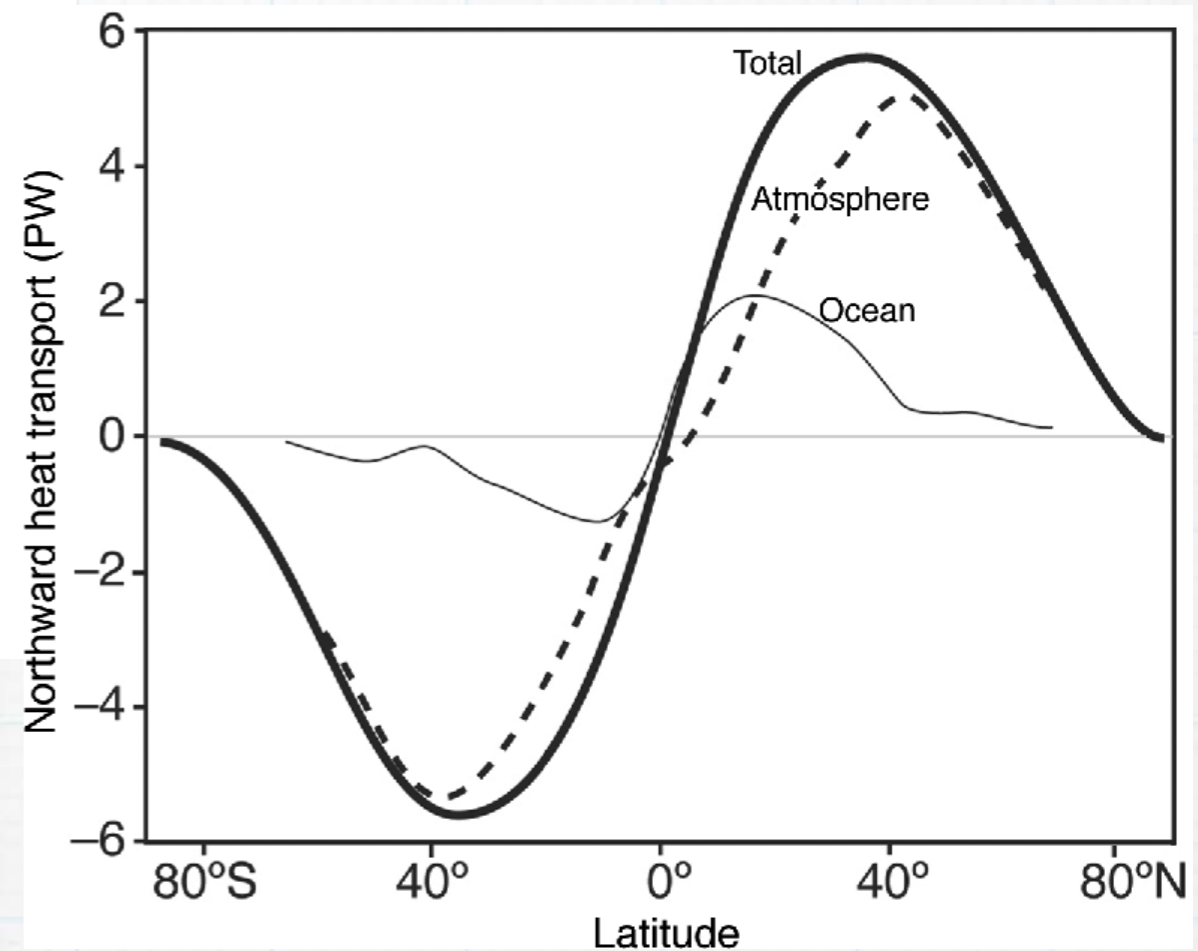
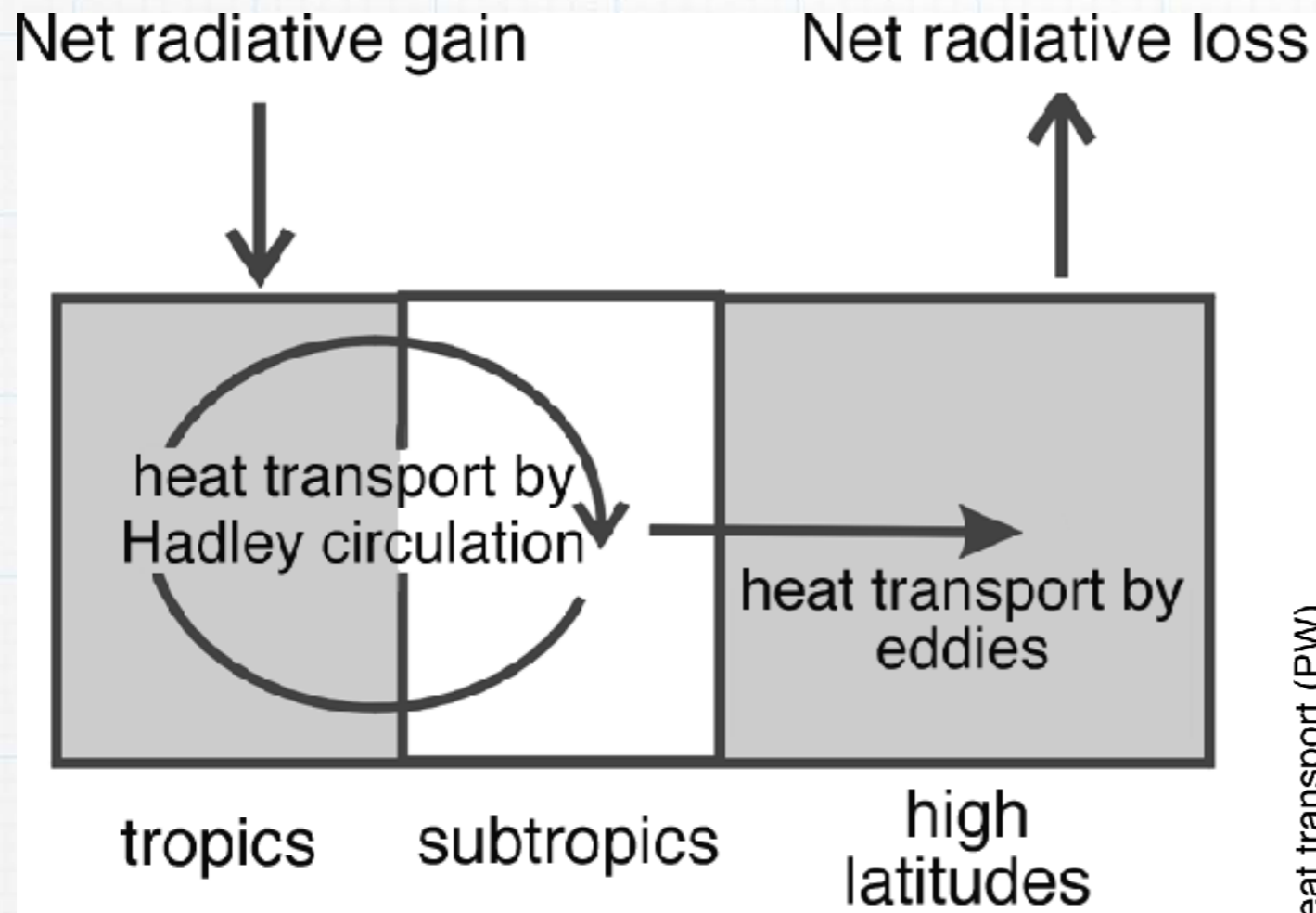
# 1. Energy transport, extratropics

- In the extratropics where the mean circulation is weak, the greater part of the transport is done by eddies.
- We saw that poleward/equatorward motions occur at almost the same altitude. → the vertical structure of the heat transport is not dominant.

- The heat transport,  $\int_0^{\infty} \rho v c_p T dz$  is positive because the poleward winds are associated with higher temperature.

- The total energy transport in the midlatitude is poleward.

# 1. Energy transport

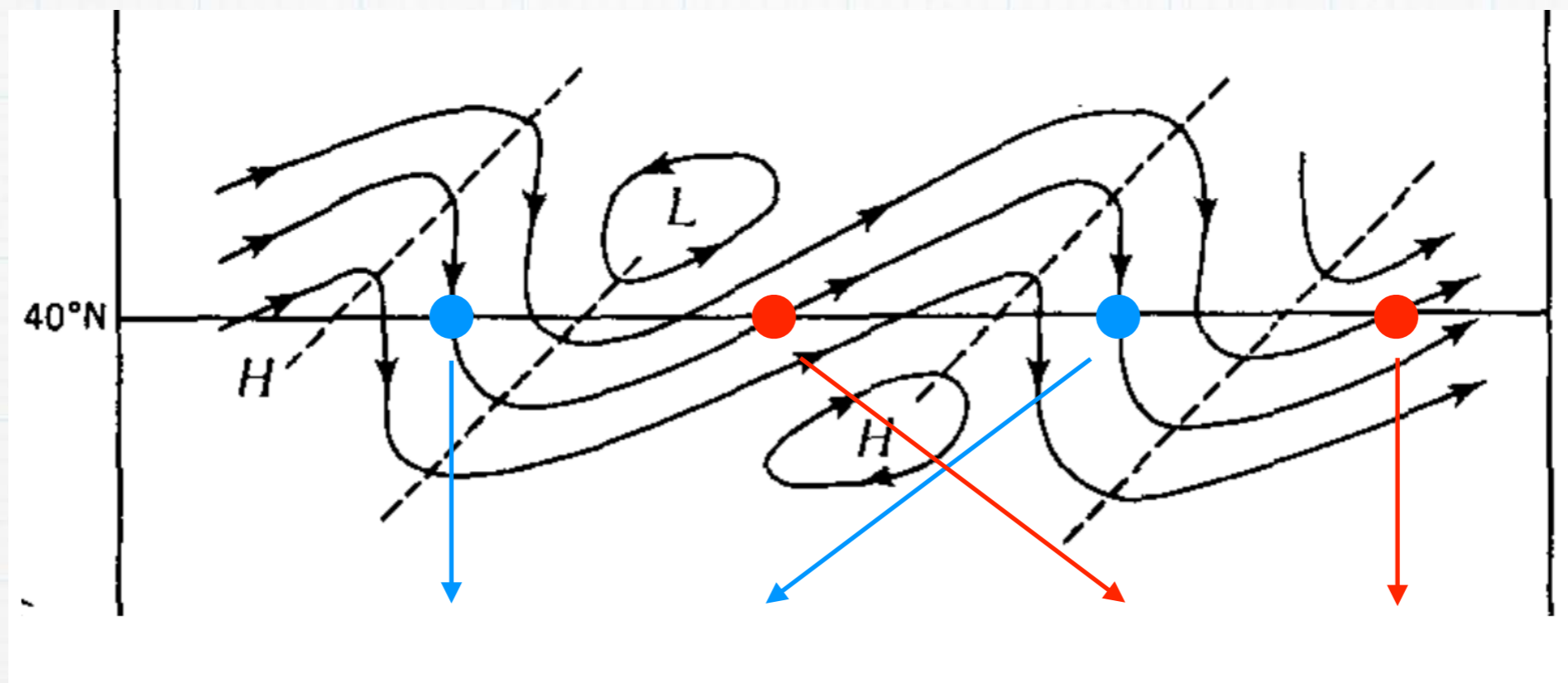


## 2. Momentum transport, tropics

- Upper branch transports westerly angular momentum poleward.
- Lower branch transport easterly angular momentum equatorward.
- Because of the friction, the momentum transport in the lower branch is weaker than the upper branch.
- The Hadley cell does a poleward transport of westerly angular momentum.

## 2. Momentum transport, extratropics

- Eddies in the extratropics also transport westerly momentum to poleward, but how?
- The meridional momentum transport  
 $= v(\Omega r^2 + ur) = v\Omega r^2 + ruv$

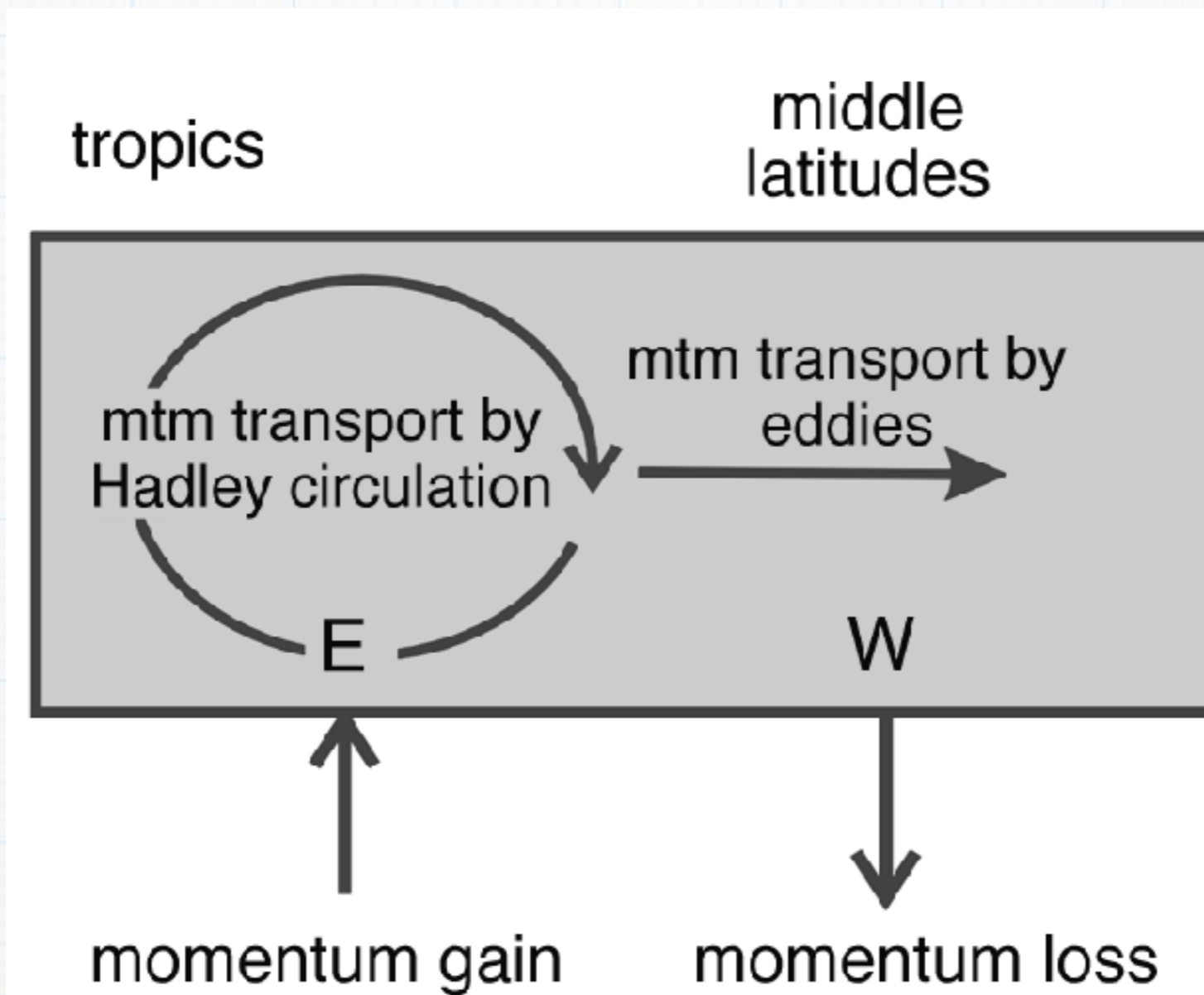


$$u \sim 0, v < 0 \rightarrow uv \sim 0$$

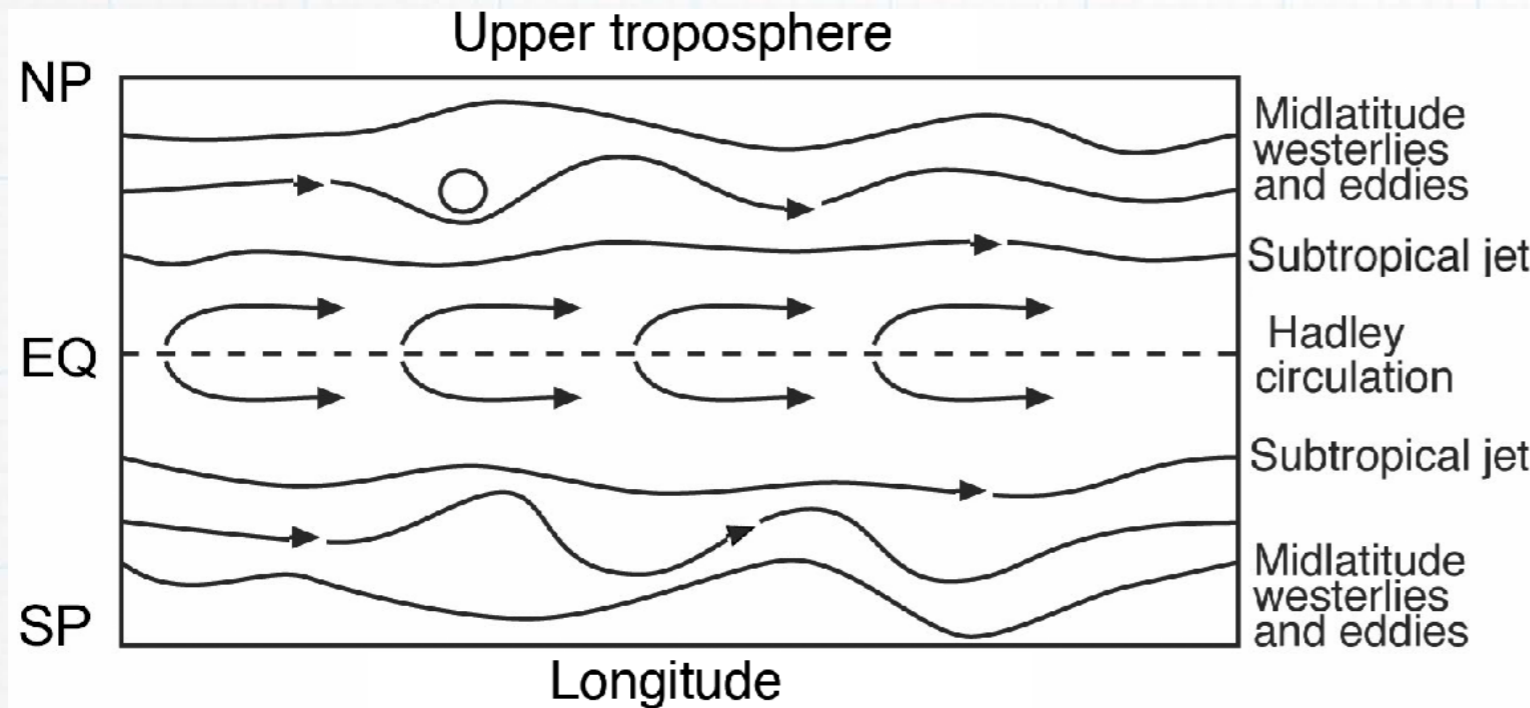
$$u > 0, v < 0 \rightarrow uv > 0$$



## 2. Momentum transport, extratropics

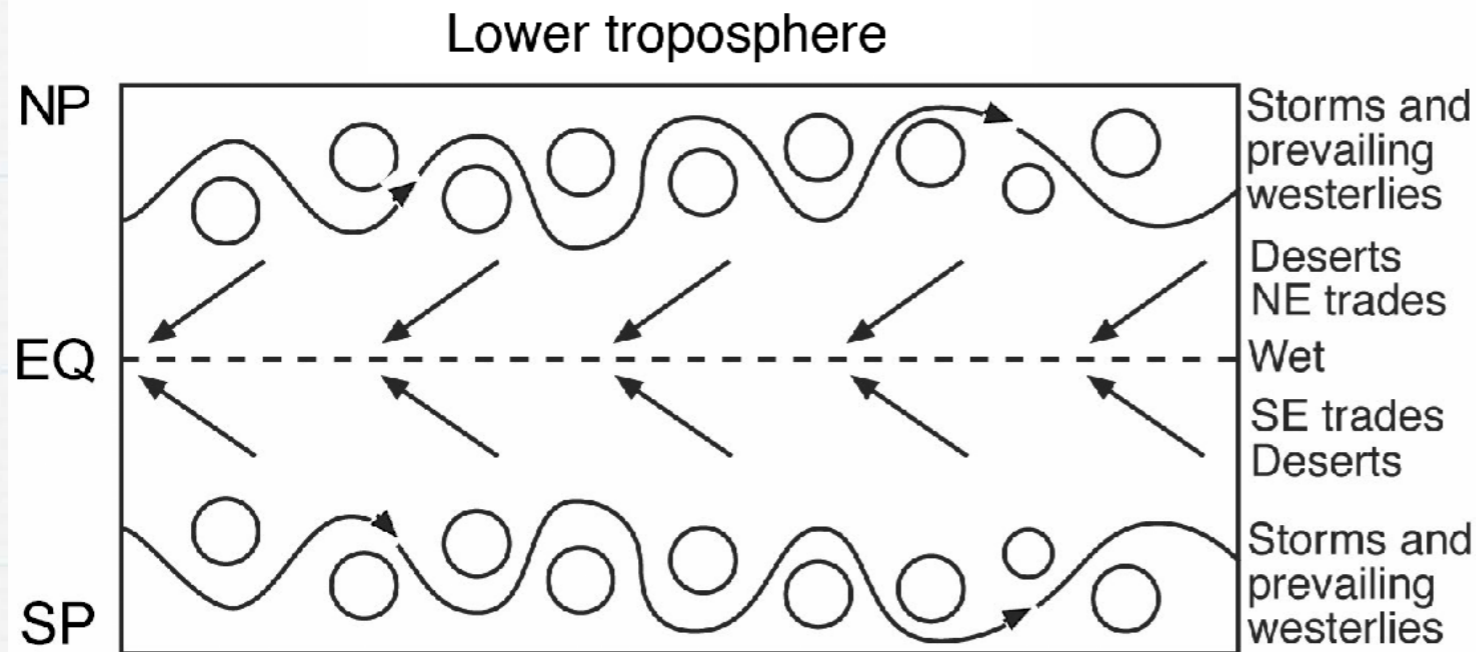


# 3. Latitudinal variations of climate



**@ tropics : convergence and upward motion → intense rainfall**

**@ midlatitude : sinking and warming → desert belt**



**@ midlatitude : eddies that go around the globe → control the weather patterns**